

A Linear Programming – Linear Assignment Approach for the Protein Morphing Problem

Sanghyun Park and Mihai Anitescu

Mathematics and Computer Science Division

Argonne National Laboratory

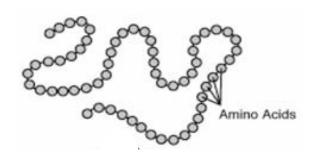


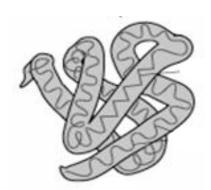
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Paradigm of Molecular Biology

Protein folding

Sequence ------ Structure ------- Function







- Function is emergent.
- The structure of the protein is key to its function.
- •Different structures = different functions

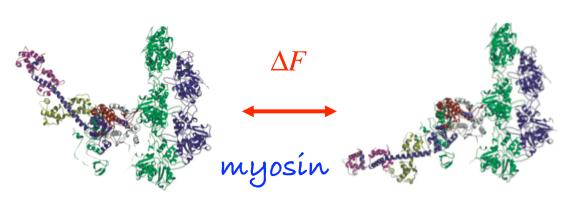
myoglobín

High concentrations of myoglobin in muscle cells allow organisms to hold their breaths longer. — Wikipedia

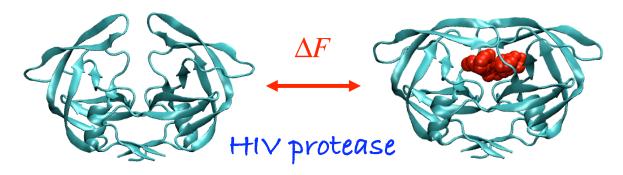


Conformational Change is Important ...

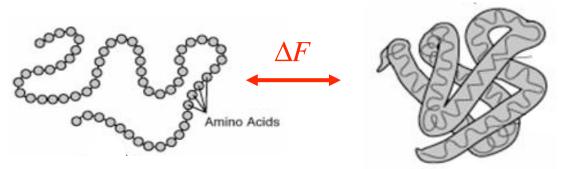
Domain motion essential to expose to binding.



Ligand binding (Drug discovery)



Protein folding

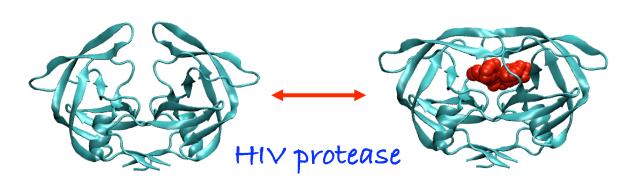




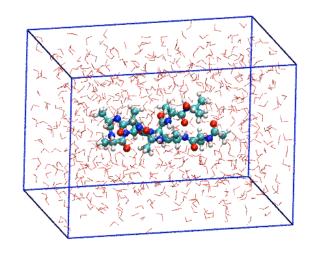
Why are free energy variations important.

- ☐ They tell whether state A is favored over state B.
- For instance, is the configuration that has the drug molecule in it favored over the one that doesn't (i.e. will this drug bind in the prescribed position at a given temperature)?
- \square The key quantity that quantifies the relative odds is $\exp\left(\frac{-\Delta E}{k_B T}\right)$

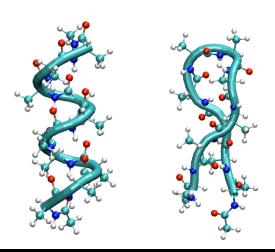
Ligand binding (Drug discovery)



Conformational Free Energy



deca-alanine



$$e^{-\beta F(\mathbf{X})} = \int d\mathbf{Y} e^{-\beta U(\mathbf{X}, \mathbf{Y})}$$

X: protein coordinates

Y: water coordinates

U(X,Y): potential energy function

F(X): conformational free energy

$$\beta = \frac{1}{k_B T}$$

$$\Delta F = F(B) - F(A)$$

Computations of F(X) don't work.



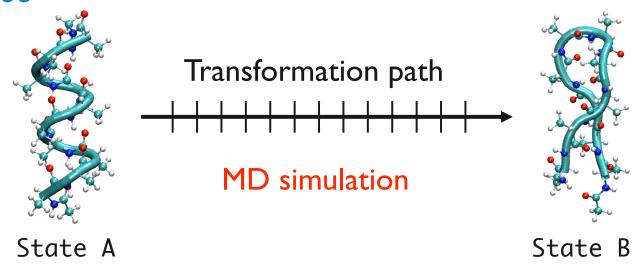
Compute ΔF directly.

Why direct computations of the free energy do not work

- Calculations of the free energy are enormously expensive.
- We have to average out the water coordinates.
- In a box, there are a few thousand water molecules, that is a few thousand water coordinates and that is for a small protein !!!
- One must estimate nasty quadrature over a 1000 dimensional space (for fixed X, the energy function has lots of minima).
- •This difficulty is known as the curse of dimensionality: The fact that in excessively large dimensions the sample density decreases exponentially.
- This is thus, a very hard problem



Free Energy Perturbation --- computing the free energy difference



Zwanzig (1954)

$$e^{-\beta\Delta F} = \left\langle e^{-\beta[U(\mathbf{A},\mathbf{Y})-U(\mathbf{B},\mathbf{Y})]} \right\rangle_{\mathbf{A}}$$

Bennett (1976): bi-directional

Jarzynski (1997): non-equilibrium

Computational cost = days x hundreds Ideal for massive parallelization





Computing free energy variations

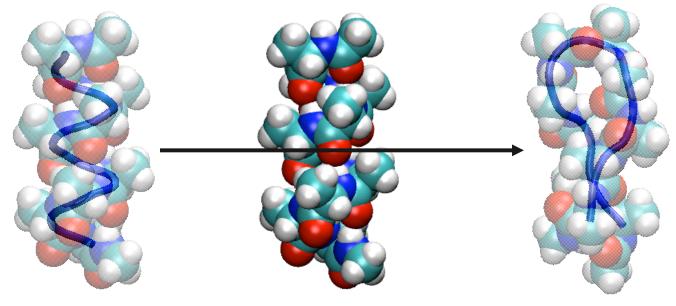
- I have just said that computing free energies is hard, so how is this possible?
- The trick is to find a "good" path in the phase space.
- Then divide the path in small segments



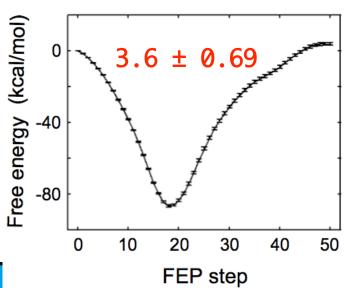
- If the path is good, the free energy difference including its variance can be computed *relatively* easily by using some version of importance sampling
- So we have transformed the problem into the one of finding a good path in phase space.



One such path: Direct Morphing

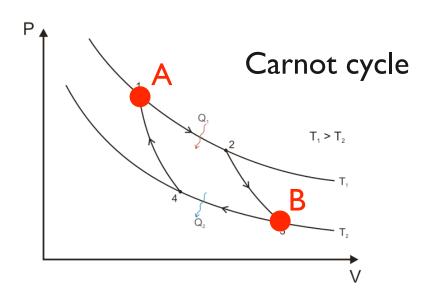


$$\mathbf{x}_n = (1 - \lambda) \, \mathbf{a}_n + \lambda \, \mathbf{b}_n$$
$$0 \le \lambda \le 1$$



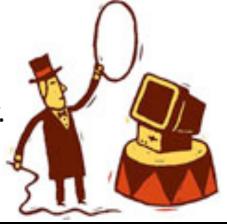


Thermodynamic Cycle



 ΔF depends only on the end states, not on the path.

Computer simulations are not bound by reality.





Computing Free Energy Between 2 States,

- We sample the distribution of states attached to each potential function "conformation" by using molecular dynamics.
- E.g., we start a molecular dynamics calculation with the potential function for fixed protein atoms, but moving water atoms, until the simulation "relaxes" and the system can be assumed ergodic.
- With the samples we create a free energy estimate using Bennett's acceptance ration method -- BAR



Bennet's Acceptance Ratio Method

We use BAR to compute the free-energy differences between neighboring states. A free-energy computation using BAR between two states proceeds as follows. A set of microstates $\{\mathbf{R}_1, \dots, \mathbf{R}_{L_1}\}$ is sampled from state 1 with potential energy function $U_1(\mathbf{R})$, and another set of microstates $\{\mathbf{R}_{L_1+1}, \dots, \mathbf{R}_{L_1+L_2}\}$ is sampled from state 2 with $U_2(\mathbf{R})$. In the present case, a microstate is a collection of protein and water coordinates, $\mathbf{R} = (\mathbf{X}, \mathbf{Y})$, except for the morphing procedure where $\mathbf{R} = \mathbf{Y}$. The free-energy difference $\Delta F := F_2 - F_1$ is then obtained by solving

$$e^{\beta \Delta F} = \sum_{l=1}^{L_1 + L_2} \left[L_1 e^{-\beta \Delta F} + L_2 e^{-\beta \Delta U(\mathbf{R}_l)} \right]^{-1}, \tag{12}$$

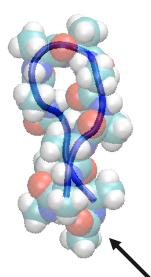
where $\Delta U = U_2 - U_1$.

- •For the equation to be nonsingular, we need the second term in the sum to be significant good overlap between ensembles.
- Energy difference is small: good transformation path AND broken down in pieces.

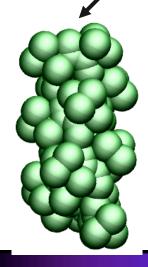


Morphing for Dummies





Alchemy



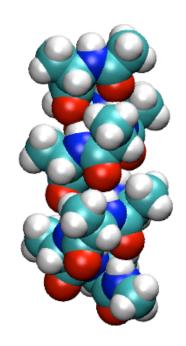
$$\mathbf{x}_n = (1 - \lambda) \,\mathbf{a}_n + \lambda \,\mathbf{b}_{\sigma(n)}$$

$$0 \le \lambda \le 1$$

Need to find a good path!

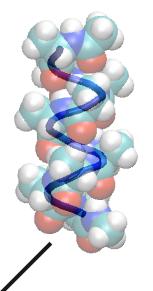
How to find a good path in phase space

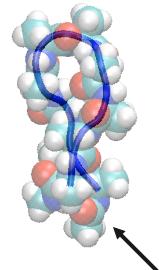
- But note that mapping the particles to the same ones from the linear structure may result sometimes in enormous traveled paths for some of them.
- And we are not bound by having each element of the trajectory feasible in the sense of it corresponding to a real compound. Such paths are very hard to find. (R. Elber, Curr. Opin. Struct. Biol. 15, 151 2005).
- Therefore, we look for different perturbations which have a chance of resulting in smaller per unit energy variations.
- •What if we actually change the atoms themselves? This will allow us to make smaller steps in energy steps at the morphing step.



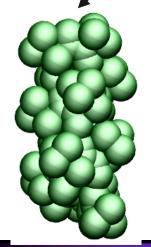


Morphing for Dummies





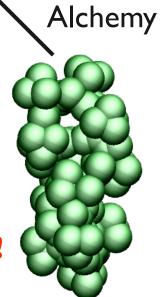
Alchemy



$$\mathbf{x}_n = (1 - \lambda) \,\mathbf{a}_n + \lambda \,\mathbf{b}_{\sigma(n)}$$

$$0 \le \lambda \le 1$$

Minimize the distances!



Least-Squares Morphing Problem

$$\min_{\sigma \in \Pi_N} \sqrt{\frac{1}{N} \sum_{n=1}^N \left\| \mathbf{a}_n - \mathbf{b}_{\sigma(n)} \right\|^2} \qquad \Pi_N = \{ N - \text{permutations} \}$$

$$\Pi_N = \{ N - \text{permutations} \}$$

$$\sum_{n=1}^{N} \|\mathbf{a}_{n} - \mathbf{b}_{\sigma(n)}\|^{2} = \sum_{n=1}^{N} \|\mathbf{a}_{n}\|^{2} + \sum_{n=1}^{N} \|\mathbf{b}_{\sigma(n)}\|^{2} - 2\sum_{n=1}^{N} \mathbf{a}_{n} \cdot \mathbf{b}_{\sigma(n)}$$

$$\max_{\sigma \in \Pi_N} \sum_{n=1}^N \mathbf{a}_n \cdot \mathbf{b}_{\sigma(n)}$$

$$P_{ij}(\sigma) = \begin{cases} 1, & j = \sigma(i) \\ 0, & \text{otherwise} \end{cases} \quad \mathbf{A} = (\mathbf{a}_1 \cdots \mathbf{a}_N) \quad \mathbf{B} = (\mathbf{b}_1 \cdots \mathbf{b}_N)$$

$$\mathbf{A} = \left(\mathbf{a}_1 \cdots \mathbf{a}_N\right) \qquad \mathbf{B} = \left(\mathbf{b}_1 \cdots \mathbf{a}_N\right)$$

$$\max_{\sigma \in \Pi_N} \operatorname{Tr} \left[\mathbf{A} \, \mathbf{P}(\sigma) \, \mathbf{B}^T \right]$$

Linear-Programming Solution

Original problem - Combinatorial search

P1:
$$\max_{\mathbf{P} \in \Omega_N} \operatorname{Tr} \left[\mathbf{A} \mathbf{P} \mathbf{B}^T \right]$$

P1:
$$\max_{\mathbf{P} \in \Omega_N} \text{Tr}[\mathbf{A} \mathbf{P} \mathbf{B}^T]$$
 $\Omega_N = \{N \times N \text{ permutation matrices}\}$



Birkoff's theorem: $\Omega_N = \{ \text{Vertices of } \Gamma_N \}$

Fundamental theorem of LP:

Solution of P2 \in {Vertices of Γ_N }

Relaxed problem - Linear programming

P2:
$$\max_{\mathbf{W} \in \Gamma_N} \operatorname{Tr} \left[\mathbf{A} \mathbf{W} \mathbf{B}^T \right] \qquad \Gamma_N = \left\{ N \times N \text{ bistochastic matrices} \right\}$$

$$\Gamma_N = \{N \times N \text{ bistochastic matrices}\}$$

$$W_{ij} \ge 0 \qquad \sum_{i} W_{ij} = 1 \qquad \sum_{j} W_{ij} = 1$$

Another way to look at it – it is a linear assignment problem!

Define
$$A = \{a_j^i\}_{i=1,p; j=1,N}$$
 $B = \{b_j^i\}_{i=1,p; j=1,N}$

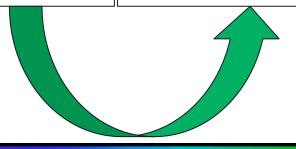
$$\max \sum_{i=1}^{N} \sum_{j=1}^{N} w_{ij} \left\langle a_{i}, b_{j} \right\rangle$$

$$\sum_{j=1}^{N} w_{ij} = 1, i = 1, 2, ..., N; \sum_{i=1}^{N} w_{ij} = 1, j = 1, 2, ..., N$$

$$\left\{ w_{ij} \right\}_{i,j=1,2...,N} \in F_{N}, w_{ij} \in \{0,1\}, i, j = 1, 2, ..., N$$

$$\max \sum_{i=1}^{N} \sum_{j=1}^{N} w_{ij} \left\langle a_{i}, b_{j} \right\rangle \\ \sum_{i=1}^{N} w_{ij} = 1, i = 1, 2, \dots, N; \sum_{i=1}^{N} w_{ij} = 1, j = 1, 2, \dots, N \\ \left\{ w_{ij} \right\}_{i, j = 1, 2, \dots, N} \in F_{N}, w_{ij} \in \{0, 1\}, i, j = 1, 2, \dots, N \\ w_{ij} \ge 0, i, j = 1, 2, \dots, N$$

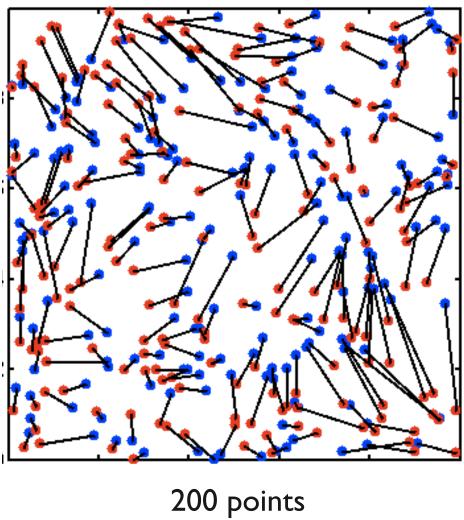
It is a linear assignment problem!



But we had to identify this in the LS formulation!!

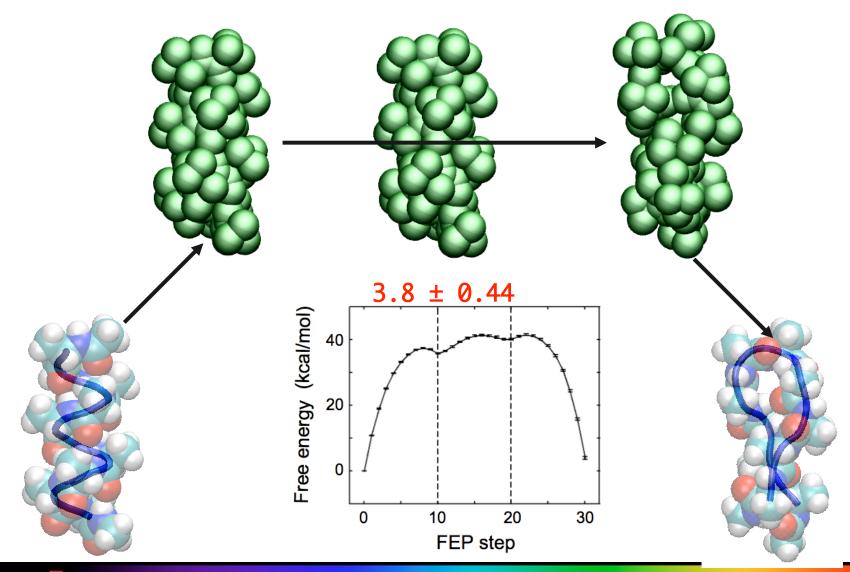


Least-Squares Permutation





Least-Squares Morphing

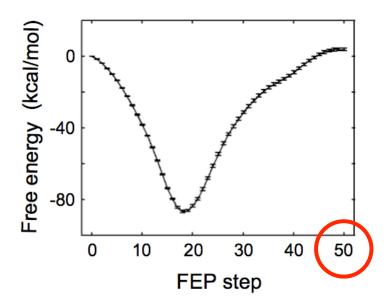




Direct vs. Least-Squares Morphing

Direct

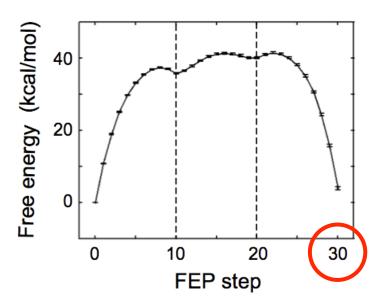
RMS distance = 8.4 Å



 3.6 ± 0.69

Least-squares

RMS distance = 2.1 Å



 3.8 ± 0.44



Discussion of the results.

- Each one of the steps in the molecular dynamics simulation is done with NAMD.
- NAMD is enormously expensive. One free energy perturbation step (FEP) takes 20 CPU hours for the deca-alanine.
- In this case, dummying the atoms takes 10 FEP, our least-squares morphing takes 10 FEPs, and the un-dummying of the atoms takes another 10 FEPs. Compare with 50 FEP steps for the original step. We save 600 CPU hours. (Morphing with LP takes 1-2 seconds).
- ■We solve 2 linear programming linear assignment problems. There are better ways to do linear assignments, but, give the small computational cost, it is not worth to do it.
- But, more importantly, we can compute a more accurate path
 0.44 versus 0.69 kcal/mol.



About NAMD- Molecular Dynamics Software

• Why it is difficult: very expensive potential – CHARM 22.

$$U(\vec{R}) = \sum_{\text{bonds}} K_{\text{b}} (b - b_{0})^{2} + \sum_{\text{UB}} K_{\text{UB}} (S - S_{0})^{2} + \sum_{\text{angle}} K_{\theta} (\theta - \theta_{0})^{2} + \sum_{\text{dihedrals}} K_{\chi} (1 + \cos(n\chi - \delta)) + \sum_{\text{impropers}} K_{\text{imp}} (\varphi - \varphi_{0})^{2} + \sum_{\text{impropers}} K_{\text{imp}} (\varphi - \varphi_{0})^{2} + \sum_{\text{impropers}} \left[\left(\frac{R_{\min_{ij}}}{r_{ij}} \right)^{12} - \left(\frac{R_{\min_{ij}}}{r_{ij}} \right)^{6} \right] + \frac{q_{i}q_{j}}{\epsilon_{1}r_{ij}}$$

- •Simulations done at constant temperature, using the Langevin thermostat and Langevin-piston barostat.
- •Time step: Ifs, it is run for Ins (1000 steps!), the trajectories sampled at 100fs are used as samples for estimating the integral.

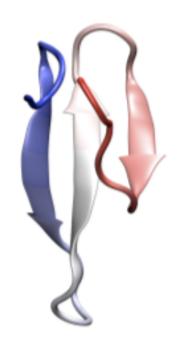


Why computing a path of small error is not trivial

- We note that getting a good path is still a matter of heuristics.
- We are interested in the overall error, not just the asymptotic error estimate for one segment, which may have the usual Monte Carlo behavior.
- Therefore it is not clear how the estimate behaves with more segments – therefore the cost of reducing the error for the original approach to the level we have obtained is hard to fathom.
- We have to some extend added a new capability to molecular dynamics.

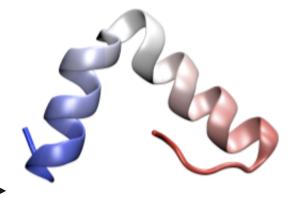


WW Domain



Direct morphing

RMS distance = 11.3 Å 100 FEP steps 12.9 ± 3.2 kcal/mol



Least-squares morphing

RMS distance = 3.4 Å 50 + 30 FEP steps $13.3 \pm 1.1 \text{ kcal/mol}$

To our knowledge, this is the first time the WW domain protein has been computed at all with this low of an error estimate.



Conclusion

- Morphing can result in much sharper estimates of free energy differences between different conformations.
- •We have shown that least-square morphing obtains an excellent free energy perturbation path.
- •We have shown that the path can be obtained in polynomial time, by using linear programming linear assignment.
- •We have obtained 100s of CPU hour computational time savings, with much more accurate FE difference estimates.



What can applied math do for FEC?

